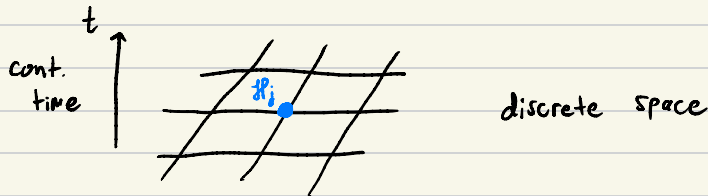


# LSM anomalies of modulated symmetries

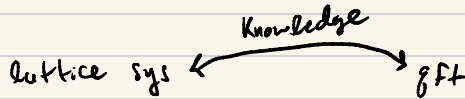
To appear tonight, w Danny Bulmash

→ Fundamental Part of a quantum theory's DNA  
 Anomalies are powerful (constrain phases, dualities, inflow)  
 - Exist on lattice and in QFT, with or without fermions

Here: Anomalies of symmetries in lattice systems



Lattice anomaly Knowledge  $\ll$  QFT anomaly Knowledge  
 (def of sym + anom on lat?)  
 (lat. and cont. form symbiotic relation.)



Morally like 't Hooft anom

Lieb-Schultz-Mattis (LSM) anomalies := obstructions to trivial symmetry-preserving gapped phases with crystalline symmetries

Eimer's talk

- original LSM anomaly: Spin 1/2 chain w  $G = so(3) \times \mathbb{Z}$  Sym.

→  $so(3)$  Sym op  $\prod_i e^{i \frac{\theta}{2} \hat{n}_i \cdot \vec{\sigma}_i}$  Spin rotation ↙ ↘ Lattice transl.

↓  
proj  $so(3)$  rep.

- Best understood for non-gen Sym with group

$$G = G_{int} \times G_{cr}$$

internal Sym group ↙ ↘ crystalline Sym group

Ho Tat's talk

- More generally

$$\mathbb{1} \rightarrow G_{\text{int}} \rightarrow G \rightarrow G_{\text{cr}} \rightarrow \mathbb{1}$$

When extension splits,  $G = G_{\text{int}} \rtimes G_{\text{cr}}$

-  $G_{\text{cr}}$  action on  $G_{\text{int}}$

$$\begin{aligned} \rho: G_{\text{cr}} &\longrightarrow \text{Aut}(G_{\text{int}}) \\ g_{\text{cr}} &\longmapsto \rho_{g_{\text{cr}}} \end{aligned}$$

- when  $\rho$  is nontrivial  $\Rightarrow$  Modulated Sym

example: dipole sym in 1+1D lattice system

- Charge operators:  $Q = \sum_j q_j$  and  $P = \sum_j j q_j$

$$TPT^\dagger = \sum_j j q_{j+1} = P + Q \Rightarrow (U(1) \times U(1)) \rtimes \mathbb{Z} \text{ sym.}$$

Modulated syms are interesting:

- SSB: goldstones in 1+1D

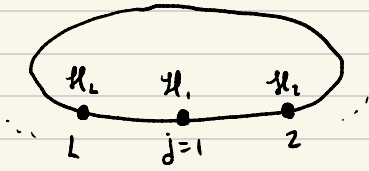
- Kramers-Wannier sym: Non-inv spatial reflections

- Gauge theory: Fractons and UV/IR mixing

This talk: LSM anomalies of modulated sym

- Restrict to bosonic systems in 1+1D with finite  $G_{\text{int}}$  and  $G_{\text{cr}}$  sym is lattice transl.

Set up: Spatial lattice of  $L$  sites on a ring



$$\mathcal{H} = \bigotimes_{j=1}^L \mathcal{H}_j$$

finite dim  $\mathcal{H}_j \cong \mathcal{H}_{\text{site}}$

$G_{\text{int}} \times \mathbb{Z}_L$  Sym ops

$$U_g = \prod_{j=1}^L U_g^{(j)} \quad \left( \begin{array}{l} \text{acts on } \mathcal{H}_j \\ (g \in G_{\text{int}}) \end{array} \right)$$

$$T^n \quad (n \in \mathbb{Z}_L)$$

With

$$U_g^{(i)} U_h^{(j)} U_{gh}^{(k)\dagger} = v_j(g, h) \in U(1) \quad (1)$$

$$T^L = 1$$

$$T U_g^{(j)} T^\dagger = U_{p_i(g)}^{(j-1)} \quad (2)$$

- From (2),  $T U_g T^\dagger = U_{p_i(g)} \quad (p: \mathbb{Z}_L \rightarrow \text{Aut}(G_{\text{int}}))$

onsite  
proj rep  $\rightarrow$

LSM anomalies arise from  $[v_j] \in \mathcal{H}^2(G_{\text{int}}, U(1))$

- conjugate (1) by  $T: v_j(g, h) = v_{j-1}(p_i(g), p_i(h))$

$$\Rightarrow [v_j] = p_i^*[v] \quad (\text{translation covariance})$$

When does  $[v]$  lead to an LSM anomaly?

Consider

$$\mathcal{H}_j = V \otimes W, \quad U_g = \prod_{j=1}^L V_g^{(j)} \otimes W_g^{(j)}, \quad T = T_V T_W$$

proj repr  $[\gamma_j]$    $[\mu_j]$

-  $G_{int} \rtimes \mathbb{Z}_L$  sym:  $T_{V,W}$  act as (2)

$$\Rightarrow [\gamma_j] = \rho_j^* [\gamma] \quad [\mu_j] = \rho_j^* [\mu]$$

- LSM anomaly-free:  $[v_j] = [\gamma_j] + [\mu_j] = 0$

Rewrite 
$$U_g = \prod_{j=1}^L W_g^{(j)} \otimes V_g^{(j+1)}$$

- Now  $[v_j] = [\gamma_{j+1}] + [\mu_j] \Rightarrow [v] = (\rho_1^* - 1)[\gamma]$

$\therefore G_{int} \rtimes \mathbb{Z}_L$  is LSM anomaly-free if  $[v] \in \text{Im}(\rho_1^* - 1)$

can show: LSM anomaly free iff  $[v] \in \text{Im}(\rho_1^* - 1)$

$P$	$[v]$	LSM anomaly
trivial or nontrivial	$= 0$	No
trivial	$\neq 0$	Yes
nontrivial	$\neq 0$	Maybe

This should be surprising!

Classification:  $H^2(G_{\text{int}}, U(1)) / \text{Im}(p_1^* - 1)$

- For infinite lattice,  $G = G_{\text{int}} \rtimes \mathbb{Z}$

$$\frac{H^2(G_{\text{int}}, U(1))}{\text{Im}(p_1^* - 1)} \cong H_{p_1^*}^1(\mathbb{Z}, H^2(G_{\text{int}}, U(1))) \subset H^3(G, U(1))$$

↑  
nontrivial  $\mathbb{Z}$ -module

Lyndon-Hochschild-Serre.

### Example: Exponential Sym

$G_{\text{int}} = \mathbb{Z}_N \times \mathbb{Z}_N$  w  $(g_1, g_2) \in \mathbb{Z}_N \times \mathbb{Z}_N$  and additive group law.

$$p_n(g_1, g_2) = (a^n g_1, b^n g_2)$$

→ Integers  $a, b$  coprime to  $N$

→ can show  $p_1^*[v] = ab[v] \quad \forall [v] \in H^2(G_{\text{int}}, U(1)) = \mathbb{Z}_N$ .

let  $[v] = K[\alpha]$  with  $[\alpha]$  a generator of  $H^2(G_{\text{int}}, U(1))$  with 2-cocycle  $\alpha(g, h) = \exp\left[\frac{2\pi i}{N} g_1 h_2\right]$

→  $[v] \in \text{Im}(p_1^* - 1)$  iff  $\gcd(ab-1, N) \mid K$ .

↳ LSM anomaly-free cond

Case	LSM anomaly for $[v] \neq 0$ ?
$ab = 1 \pmod{N}$	Always
$\gcd(ab-1, N) = 1$	Never
$\gcd(K, N) = 1$	iff $\gcd(ab-1, N) \neq 1$

## Outlook

LSM anomalies of modulated sym w lattice transl in 1+1D

→ LSM anomaly existence depends on onsite proj rep and homomorphism  $\rho$ .

What else is in arXiv: 2602.XXXXX

1) Examples of 1+1D stabilizer code models

- SPT-LSM theorems

2) LSM anomalies of Gint XI pm

3) Beyond 1+1D

- Stratified anomalies: generalization of onsite proj reps

4) Examples in 2+1D.